# Effect of Magnetic Field and Temperature on the Exchange Bias in Phase-Separated Nd<sub>1-x</sub>Sr<sub>x</sub>CoO<sub>3</sub>(x=0.10, 0.15)

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**Abstract:** The exchange bias phenomena of phase-separated  $Nd_{1-x}Sr_xCoO_3(x=0.10, 0.15)$  samples were systematically investigated in this paper. The samples were prepared using conventional solid state reaction method. When the  $Nd_{1-x}Sr_xCoO_3$  samples cooled down in magnetic field below freezing temperature, the hysteresis loops shifted along the magnetic field axis. Moreover, exchange bias of  $Nd_{1-x}Sr_xCoO_3$  is strongly dependent on the field and the temperature. The influence of magnetic field on the relative ratio of the coexisting phases may be responsible for these behaviors. Therefore, our study confirmed that in phase-separated system, the exchange coupling at the interface between the ferromagnetism clusters and the spin glass regions may induce interfacial exchange anisotropy.

Keywords: Phase separation, exchange bias, exchange coupling.

## **1. INTRODUCTION**

When samples with a ferromagnetic-antiferromagnetic (FM-AFM) interface cool down in the magnetic field from above Neel temperature, hysteresis loops shift towards the magnetic field axis when the hysteresis loops are measured under the Neel temperature, namely the exchange bias. Exchange bias is induced by unidirectional anisotropy caused by exchange-coupling at the FM-AFM interface. In 1956, Meiklejion et al. found exchange bias by studying Co particle shrouded by CoO. Ever since, people have found exchange bias at FM-AFM /Ferri-FM/Ferri-AFM interfaces [1-6]. Exchange bias exists in mono-crystal and poly-crystal samples and thin film samples [7-15], suggesting that exchange bias is an internal trait of these systems, but not related to synthesis methods. In the recent years, exchange bias has also been found in the spin glass. Exchangecoupling at ferromagnetic and spin glass interface induces exchange bias [16-18].

This thesis studied the exchange bias of cobalt oxide Nd<sub>1-</sub> $_x$ Sr<sub>x</sub>CoO<sub>3</sub> (*x*=0.10, 0.15). In research, it is considered that cobalt oxides are composed of ferromagnetic cluster, non-ferromagnetic matrix and spin glass, and that spin glass exists between ferromagnetic cluster and non-ferromagnetic matrix, forming an interface around ferromagnetic cluster [19-21]. Under such circumstance, phase separation of Nd<sub>1-</sub> $_x$ Sr<sub>x</sub>CoO<sub>3</sub> causes the formation of a ferromagnetic /spin glass interface. Hence, we studied the hysteresis loops of Nd<sub>1-</sub> $_x$ Sr<sub>x</sub>CoO<sub>3</sub>(*x*=0.10, 0.15) when the field cooled down.

Results show that exchange bias is strongly dependent on the temperature and the magnetic field.

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### 2. EXPERIMENT

Poly-crystal samples  $Nd_{1-x}Sr_xCoO_3(x=0.10, 0.15)$  were synthesized with standard solid-state reaction method. Firstly, highly purified  $Nd_2O_3$  was pre-sintered for four hours at 800 °C. Subsequently, a stoichiometric mixture of  $Nd_2O_3$ ,  $SrCO_3$  and  $Co_2O_3$  powder was acquired; it was then presintered for 48 hours at 1 000 °C after grinding and tablet compressing. Next, regrinding and re-tablet compressing occurred, and the sintering temperatures were set as 1 100 °C and 1 200 °C respectively, while the sintering time was set at 48 hours; the sample synthesis was completed until the temperature dropped down gradually to room temperature for 12 hours.



**Fig. (1).** XRD spectrogram of poly-crystal sample  $Nd_{1-x}Sr_xCoO_3$  (*x*=0.10, 0.15) at room temperature.

X-ray diffraction (XRD) (Fig. 1) power pattern was collected using the Bede D<sup>1</sup> XRD spectrometer with Nifiltered Cu  $K\alpha$  radiation. The room temperature XRD pattern showed that the sample is single phase with orthorhombic structure. Magnetic testing adopts the physical properties measurement system (PPMS) of the Quantum Design.

### **3. RESULTS AND DISCUSSION**

For the condition of field-cooled (FC), hysteresis tests were conducted on samples (x=0.10) at 10 K and 25 K to analyze the causes of exchange bias. For each measurement, the field is always 5 KOe when the temperature drops from 300 K to the testing temperature.

Fig. (2) shows the *M*-H curve (x=0.10) at 10 K. The illustration is an enlarged view of the center of M-H curve. It can be observed from the figure that samples are characterized by obvious exchange bias at 10 K:  $H_{\rm E}$ =50 Oe,  $M_{\rm E}$ =0.037 emu/g. The vertical movement in the center of the hysteresis loops is defined as  $M_{\rm E}$ , and horizontal movement as  $H_{\rm E}$ .  $M_{\rm E}=(M_{\rm up}+M_{\rm down})/2$ .  $M_{\rm up}$  and  $M_{\rm down}$  are the intersections of hysteresis loops with M;  $H_{\rm E}=-(H_{\rm right}+H_{\rm left})/2$ .  $H_{\text{right}}$  and  $H_{\text{left}}$  are the intersections of hysteresis loops with H. Fig. (3) shows the M-H curve (x=0.10) at 25 K. The illustration at the bottom right is an enlarged view of the center of M-H. It can be observed from the figure that when the temperature reaches 25 K (above the freezing temperature),  $H_{\rm E}$  and  $M_{\rm E}$  of the samples are both zero, and exchange bias disappears. But in the upper left illustration of Fig. (3), samples (x=0.15) still display exchange bias when the temperature reaches 25 K (below the freezing temperature):  $H_{\rm E}$ =283Oe,  $M_{\rm E}$ =0.52emu/g.



**Fig. (2).** The *M*-*H* curve (x=0.10) at 10 K on condition of  $H_{cool}=100$  Oe. The illustration is an enlarged view of the center of *M*-*H* curve.

Around freezing temperature, exchange bias disappears, which is very typical at FM/SG interfaces [12, 16-17]. When samples cool down on the condition of FC past  $T_{\rm C}$ , the magnetic moment of the ferromagnetic layer turns to the external magnetic field, while that of spin glass is still characterized by the disorder. When the temperature drops below the freezing temperature, due to exchange-coupling at

the ferromagnetic cluster and spin glass interface, the magnetic moment at the spin glass interface may not randomly freeze. Due to exchange-coupling, the magnetic moment of the spin glass layer is arranged in a certain direction close to the ferromagnetic cluster. When the direction of the external field is reversed, the magnetic moment in the ferromagnetic region starts to spin, while that in the SG region stays unchanged. Besides, due to exchangecoupling, the magnetic moment in the spin glass interface exerts the influence of torque on the magnetic moment of the ferromagnet, keeping the original spin direction of the ferromagnetic cluster. Since the effect of magnetic moment needs to be overcome, the magnetic field that changes the ferromagnetic is high. When the direction of the external field is changed again, due to the effect of micro-torque, the magnetic moment at a small magnetic field spins. Thus,  $H_{\rm E}$ appears referred to as the exchange bias.



**Fig. (3).** The *M*-*H* curve (x=0.10) at 25 K on condition of  $H_{cool}$ =100 Oe. The illustration at bottom right is the enlarged view of the center of *M*-*H*; the upper left illustration is the curve (x=0.15) at 25 K.

The hysteresis loops display vertical movement below the freezing temperature, indicated by  $M_{\rm E}$ . This is because there is a competition of energies at FM/SG interface. In the external magnetic field, magnetic behavior is mainly influenced by four energies, including anisotropy energy of FM cluster  $E_{\rm F}$  and Zeeman energy  $E_{\rm Z}$ , exchange energy at FM/SG interface  $E_{int}$  and anisotropy energy of SG  $E_S$  [3-4,16]. Effective Zeeman energy  $E_{\text{Zeff}} = |E_Z| - |E_F|$  is introduced. When  $E_{\text{Zeff}} < E_S$  and  $E_{\text{Zeff}} < E_{\text{int}}$ , Zeeman energy is feeble and it is not able to overcome the barrier at the interface or to reverse the magnetic moment of the spin glass. After field cooling, the magnetic moment of ferromagnetic cluster gets frozen in the direction of the external field. Thus, the hysteresis loops display biased magnetic moment in the vertical direction. For samples (x=0.1, 0.15), the fields are both 5 KOe. Zeeman energy  $E_Z$  is feeble. Thus, the bias  $M_E$ appears in the direction of the magnetic moment of hysteresis loops.



Fig. (4). the M-H curve of Nd<sub>0.85</sub>Sr<sub>0.15</sub>CoO<sub>3</sub> samples at 25 K on condition of FC.

The research reveals the relevance between exchange bias and the external field [7, 12, 16, 18]. Fig. (4) includes hysteresis loops at different fields on condition of FC (x=0.15). When the field is 5 KOe, exchange bias is apparent.  $H_{\rm E}$ =283 Oe,  $M_{\rm E}$ =0.52emu/g. When the field is 50 KOe, exchange bias disappears.

While the field increases, exchange bias is apparently lessened and disappears in high magnetic field. This is because the relative ratio of coexisting phases in the system is affected by the external field. The same can be found in the research conducted by Tang, *et al.* on  $La_{1-x}Sr_xCoO_3[16]$ , and Mao, et al. on La<sub>0.8</sub>Sr<sub>1.2</sub>CoO<sub>4</sub>[18]. On one hand, in Nd<sub>1</sub>.  $_x$ Sr<sub>x</sub>CoO<sub>3</sub>(x=0.10, 0.15), ferromagnetic cluster increases in external field while the spin glass decreases in the external field. The proportion of spin glass decreases with the increase in ferromagnetic cluster and the damage of spin glass. When the external field is high enough, relatively speaking, the effect of the magnetic moment of spin glass on that of ferromagnetic is feeble, and the ferromagnetic cluster is completely reversed in high magnetic field. Thus, the exchange anisotropy at the interface disappears and the horizontal movement  $H_{\rm E}$  of hysteresis loops disappears. This is different from the traditional exchange bias. The exchange bias is strongly dependent on the external field. This dependence suggests that when phase-separated cobalt oxide cools down in static magnetic field, the exchange effect at ferromagnetic cluster and spin glass interface induces unidirectional anisotropy. On the other hand, since Zeeman energy increases with the increase in field, under the influence of external field, frozen ferromagnetic cluster starts to spin and vertical movement  $M_{\rm E}$  of hysteresis loops decreases. When Zeeman energy increases to the extent where it can overcome the coupling of the interface, the magnetic moment of ferromagnetic is completely reversed. In the high magnetic field,  $M_{\rm E}$  disappears [19-21].

### CONCLUSION

To conclude, this thesis observed exchange bias in phaseseparated cobalt oxide  $Nd_{1-x}Sr_xCoO_3(x=0.10, 0.15)$ . Through studying the effect of exchange bias of Nd<sub>1-x</sub>Sr<sub>x</sub>CoO<sub>3</sub>, it can be observed that exchange bias in the system is strongly dependent on the magnetic field and temperature. Since different phases in Nd<sub>1-x</sub>Sr<sub>x</sub>CoO<sub>3</sub> change in different fields, when the external field is high enough, the spin glass is damaged, the ferromagnetic cluster is reversed and the exchange bias disappears. Thus, the research on the exchange bias suggests that in Nd<sub>1-x</sub>Sr<sub>x</sub>CoO<sub>3</sub>(x=0.10, 0.15), located between the ferromagnetic and non-ferromagnetic field, the spin glass forms an interface. Since there is exchange between the ferromagnetic cluster and the spin glass, unidirectional anisotropy is induced at the interface, and thus the appearance of exchange bias occurs.

### **CONFLICT OF INTEREST**

The author confirms that this article content has no conflict of interest.

### **ACKNOWLEDGEMENTS**

The present work is supported by the Special Scientific Research Program Funded of Ankang University (Program No. 2015AYPYZR05), Project from Scientific Research Fund of Shaanxi Provincial Education Department (Grant No. 2013K1115).

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#### Effect of Magnetic Field and Temperature on the Exchange Bias

#### The Open Materials Science Journal, 2015, Volume 9 59

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Revised: May 21, 2015

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Accepted: June 9, 2015

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Received: February 17, 2015