

The Transient Method and Experimental Study on Threshold Pressure Gradient of Heavy Oil in Porous Media

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Abstract: On the basis of flow characteristics and material balance principle, a one-dimensional theoretical model of heavy-oil migration in porous media was established. The model considered micro-compressible property and non-Newtonian behavior of heavy oil. A new method, the transient method, was introduced to calculate the threshold pressure gradient in porous media. The method was designed to measure the inlet pressure of porous media during experiments. Finally, the threshold pressure gradient was worked out according to the mathematical model. The threshold pressure gradient of heavy oil was more accurate by the transient method than the steady state method. The threshold pressure gradients of three oil samples in different permeable porous media were measured according to the transient method. The experimental results showed that heavy oil migration was influenced by petroleum composition and pore structure of porous media. The threshold pressure gradient gradually decreased as temperature and permeability increased in porous media.

Keywords: Heavy oil, Threshold pressure gradient, Transient method, Porous media, Inflection temperature.

1. INTRODUCTION

Heavy oil and bitumen were important hydrocarbon resources that played an increasingly great role in petroleum supply for the world [1-3]. The total amount of the heavy oil resources all over the world were 10 trillion barrels, nearly three times of the conventional oil in the world. Heavy oils could not be recovered by traditional ways due to their high viscosity and apparent flow resistance in oil reservoirs. Thermal oil recovery methods have been the most effective enhanced oil recovery technologies to recover heavy oil [4]. Thermal methods aimed at reducing oil viscosity to increase its mobility, which mainly included in-situ combustion, steam drive, cyclic steam stimulation and steam assisted gravity drainage.

Heavy oils were complex fluids that could cause a variety of difficulties during production, separation, transportation and refining [5, 6]. Heavy oil always posed a great challenge to production and transportation systems due to its high viscosity. Meanwhile, viscosity was related to API gravity, solution gas oil ratio, pressure and temperature. Many researchers explained that heavy oil viscosity usually varied dramatically during various production processes, such as a thermal or solvent injection processes [7-10].

Luo and Gu (2007) [11] carried out many experiments to study thermal flow properties of heavy oil. The results focused on the relationships between the compositions of heavy oils, in particular asphaltine contents, and their flow

properties. Some researchers observed that as temperature decreases, not only the viscosity increased but also the rheological behavior of heavy oils became non Newtonian [12-14]. The heavy oil exhibited non-Newtonian flow behavior of shear thinning, which could be best presented by the power law model [15, 5].

Generally, heavy oil contained high contents of large molecular hydrocarbon and asphaltenes but only a small amount of light hydrocarbons and easily volatile components [4, 7, 11, 13]. Therefore, heavy oil was a kind of high viscosity and large density of fluid, which difficultly flowed in porous media under natural conditions. The results from experiments and oil-field applications showed that the seepage of heavy oil was presented nonlinear in porous media [16, 17]. Heavy oil could be driven in porous media until pressure difference was more than a critical value that was called threshold pressure. The steady-state method was very difficult to accurately control flow rate of heavy oil and perfectly measured pressure difference, especially, at lower flow rate [18-21]. Meanwhile, the steady-state method needed to carry out many experiments at different flow rates.

In this article, according to seepage characteristics of heavy oil and mass balance theory, a transient method was presented to measure threshold pressure gradient of heavy oil in porous media. This method only needed to monitor pressure variation rather than flow rate of heavy oil in sand-pack. Therefore, it was superior to the steady-state method.

2. THEORY

2.1. Mathematic Model

Generally, heavy oil reservoirs presented some characteristics, such as loose sand grains, high permeability

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and large porosity [1]. A long sand-pack filled with uniform quartz grains or glass beads was employed to measure threshold pressure gradient of heavy oil in porous media. A transient model was established on the basis of unsteady seepage theory and material balance principle. The following basic assumptions were presented: (1) Porous media was incompressible and homogeneity; (2) Temperature was same in porous media; (3) Crude oil was slightly compressible. Therefore, the following equations could be established.

(1) Motion equation

$$\begin{cases} \bar{v} = -\frac{K}{\mu_o}(\text{grad}p + \lambda) & |\text{grad}p| > \lambda \\ \bar{v} = 0 & |\text{grad}p| \leq \lambda \end{cases} \quad (1)$$

(2) Conservation equation

$$\frac{\partial(\rho_o \phi)}{\partial t} + \frac{\partial(\rho_o \bar{v})}{\partial x} = 0 \quad (2)$$

(3) Equation of state

Density of heavy oil under pressure, p , was:

$$\rho_o = \rho_{osc} e^{c_o(p-p_{sc})} \quad (3)$$

(4) Boundary and initial conditions

$p|_{x=L} = p_{sc}$, corresponding to the outlet pressure of sand pack.

$\frac{\partial p}{\partial x}|_{x=0} = -\lambda$, corresponding to shutting in the inlet of sand pack, that is, no flow at the inlet.

$p|_{t=0} = p_i$, corresponding to the initial pressure of sand pack.

2.2 Solution Method

(1) Simplified equation

$$\frac{\partial(\rho_o \bar{v})}{\partial x} = -\frac{K \rho_{osc}}{\mu_o} \left\{ \frac{1}{c_o} \frac{\partial^2 [e^{c_o(p-p_{sc})}]}{\partial x^2} + \lambda \frac{\partial [e^{c_o(p-p_{sc})}]}{\partial x} \right\} \quad (4)$$

If $\bar{p} = e^{c_o(p-p_{sc})}$, then

$$\begin{cases} \frac{\partial^2 \bar{p}}{\partial x^2} + c_o \lambda \frac{\partial \bar{p}}{\partial x} = \frac{c_o \phi \mu_o}{K} \frac{\partial \bar{p}}{\partial t} \\ \bar{p}|_{t=0} = e^{c_o(p_i-p_{sc})} = \bar{p}_i \\ \bar{p}|_{x=L} = 1 \\ \left. \frac{1}{\bar{p}} \frac{\partial \bar{p}}{\partial x} \right|_{x=0} = -c_o \lambda \end{cases} \quad (5)$$

(2) Dimensionless equation

If $p_D = \frac{\bar{p}}{\bar{p}_i}$, $X_D = \frac{x}{L}$, $t_D = \frac{K}{\phi \mu_o c_o L^2} t$, $\lambda_D = L c_o \lambda$, then

$$\begin{cases} \frac{\partial^2 p_D}{\partial x_D^2} + \lambda_D \frac{\partial p_D}{\partial x_D} = \frac{\partial p_D}{\partial t_D} \\ p_D|_{t_D=0} = \bar{p}_i \\ p_D|_{x_D=1} = 1 \\ \left. \frac{1}{p_D} \frac{\partial p_D}{\partial x_D} \right|_{x_D=0} = -\lambda_D \end{cases} \quad (6)$$

(3) Solution of equations

In order to obtain the solution, we assumed that

$$p_D(x_D, t_D) = u(x_D, t_D) + w(x_D) \quad (7)$$

Therefore,

$$\begin{cases} \frac{\partial^2 u}{\partial x_D^2} + \lambda_D \frac{\partial u}{\partial x_D} + \frac{d^2 w}{dx_D^2} + \lambda_D \frac{dw}{dx_D} = \frac{\partial u}{\partial t_D} \\ u|_{t_D=0} = \bar{p}_i - w(x_D) \\ u|_{x_D=1} = 1 - w(1) \\ \left. \left(\frac{\partial u}{\partial x_D} + \lambda_D u \right) \right|_{x_D=0} = - \left. \left(\frac{dw}{dx_D} + \lambda_D w \right) \right|_{x_D=0} \end{cases} \quad (8)$$

If we assumed that $w(x_D)$ was satisfied with the ordinary differential equations as following:

$$\begin{cases} \frac{d^2 w}{dx_D^2} + \lambda_D \frac{dw}{dx_D} = 0 \\ w(1) = 1 \\ \left. \left(\frac{dw}{dx_D} + \lambda_D w \right) \right|_{x_D=0} = 0 \end{cases} \quad (9)$$

We could derive the following solution:

$$w(x_D) = e^{\lambda_D(1-x_D)} \quad (10)$$

Therefore, the equation (8) could be simplified to

$$\begin{cases} \frac{\partial^2 u}{\partial x_D^2} + \lambda_D \frac{\partial u}{\partial x_D} = \frac{\partial u}{\partial t_D} \\ u|_{t_D=0} = \bar{p}_i - e^{\lambda_D(1-x_D)} \\ u|_{x_D=1} = 0 \\ \left. \left(\frac{\partial u}{\partial x_D} + \lambda_D u \right) \right|_{x_D=0} = 0 \end{cases} \quad (11)$$

Let $u(x_D, t_D) = e^{-\frac{\lambda_D x_D}{2}} e^{C_0 t_D} v(x_D, t_D)$ and $C_0 = -\frac{\lambda_D^2}{4}$, then

$$\begin{cases} \frac{\partial^2 v}{\partial x_D^2} = \frac{\partial v}{\partial t_D} \\ v|_{t_D=0} = e^{\frac{\lambda_D x_D}{2}} [\bar{p}_i - w(x_D)] = F(x_D) \\ v|_{x_D=1} = 0 \\ \left. \left(\frac{\partial v}{\partial x_D} + \frac{\lambda_D}{2} v \right) \right|_{x_D=0} = 0 \end{cases} \quad (12)$$

If equation (12) was solved according to separation of variable by superposition, we could derive

$$\beta_m \text{ctg}(\beta_m) = \frac{\lambda_D}{2} \quad (13)$$

If we used Normalization factor of $N = \int_0^1 \sin^2 \beta_m (1-x) dx$, then

$$\frac{1}{N(\beta_m)} = \frac{8\beta_m^2 + 2\lambda_D^2}{4\beta_m^2 + \lambda_D^2 - 2\lambda_D} \quad (14)$$

The expression of p_D was worked out from the equations of (6), (11) and (12).

$$p_D = e^{-\frac{\lambda_D}{2}x_D} e^{-\frac{\lambda_D^2}{4}t_D} \left[\sum_{m=1}^{\infty} e^{-\beta_m^2 t_D} \frac{X(\beta_m, x_D)}{N(\beta_m)} \int_0^1 X(\beta_m, x_D) F(x_D) dx_D \right] + e^{\lambda_D(1-x_D)} \quad (15)$$

Where, $X_m = \sin[\beta_m(1-x_D)]$;

$$F(x_D) = e^{\frac{\lambda_D}{2}x_D} [1 - e^{\lambda_D(1-x_D)}]$$

3. EXPERIMENTS

3.1. Experimental Apparatus

The flow experiments of heavy oil in porous media were carried out using a sand-pack flooding unit, which included constant flow pump, heavy oil tank, sand-pack, constant temperature oven, pressure transducer, temperature sensor, pressure gauge, data acquisition systems and computer. A schematic of the unit was shown in Fig. (1). The sand-pack was 60 cm long with 3.8 cm in inner diameter, in which was filled with unconsolidated quartz grains. The sand-pack was equipped with three pressure transducers and three temperature sensors. In order to prevent heavy oil from plugging pipelines, one small buffer container, which was a small hollow tank, was installed between pressure tap and

pressure transducer. Heavy oil in oil tank was injected into sand-pack by a constant flow bump. Heavy oil tank, sand-pack, small buffer containers and temperature sensors were all equipped in the constant temperature oven. The data of temperatures and pressures were transferred to a computer by the data acquisition systems.

3.2. Experimental Procedure

(1) The sand-pack filled with quartz grains was first equipped in the constant temperature oven. Then the oven and the oil tank were controlled to (in) a certain temperature over 2 hours.

(2) Heavy oil was injected into sand-pack until the flow rate of heavy oil was steady in the outlet of sand-pack.

(3) The valve of the outlet was shut in but heavy oil was still injected into sand-pack by the constant flow bump until the inlet pressure was 0.5 MPa higher than reservoir original pressure. Then the valve of the inlet was shut in.

(4) The temperature of the oven was evaluated about 10 °C in order to accelerate the pressure balance of heavy oil in sand-pack.

(5) When the pressures were wholly equal to each other from the inlet to the outlet of sand-pack, the valve of the outlet was opened. Simultaneously, the inlet pressure of sand-pack was continuously recorded until the inlet pressure was steady.

(6) The dimensionless pressure of the inlet at a certain temperature could be calculated according to the equation (15). At last the threshold pressure gradient was calculated resulting from experimental results.

3.3. Validation of Transient Method

A heavy oil sample was employed to measure its threshold pressure gradient in sand-pack respectively at 30

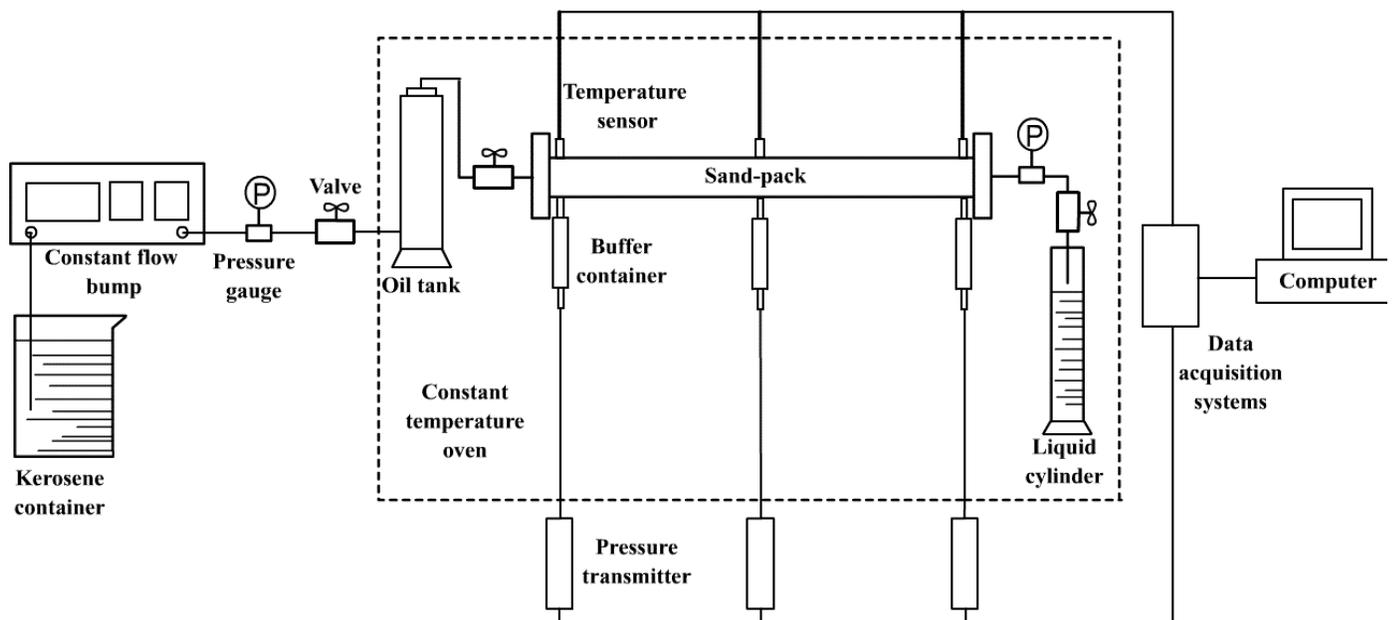


Fig. (1). The Schematic of heavy oil threshold pressure gradient experiment.

and 50 °C. Firstly, a series of experiments were carried out to establish the relationships between pressure gradient and flow rate of heavy oil in porous media at steady state. The threshold pressure gradient of this oil sample was calculated when flow rate was equal to zero according to the experimental results. The experimental parameters were listed in Table 1. The permeability of sand-pack was 2.55 μm^2 , which was measured according to Darcy’s Law. The porosity of sand-pack was 0.37.

Table 1. The Parameters of Porous Media and Oil Properties

Parameters		Value	
Sand-pack	Length (cm)	60.0	
	Diameter (cm)	3.8	
Porous media	Quartz grain	160 meshes	
	Porosity (%)	37.0	
	Permeability (μm^2)	2.55	
Heavy oil	Compressibility (1/MPa)	4.5×10^{-3}	
	Temperature	30	6110.00
	-viscosity	50	1401.44
	(°C, mPa·s)	80	237.04

The relationships between pressure gradient and flow rate at steady state were shown in Fig. (2). At experimental temperature, pressure gradient was linear with flow rate but the line did not pass through the origin of coordinate axis. It showed that heavy oil began to flow in porous media when displacement pressure difference was larger than threshold pressure difference. Threshold pressure gradient was 0.3583 MPa/m at 30 °C for experimental oil sample. It decreased to be 0.1554 MPa/m at 50 °C. Therefore, threshold pressure gradient could be decreased to zero when temperature increased to a certain value, that is, heavy oils changed to Newton fluid from visco-plastic fluid at the temperature what was called inflection temperature [5, 9].

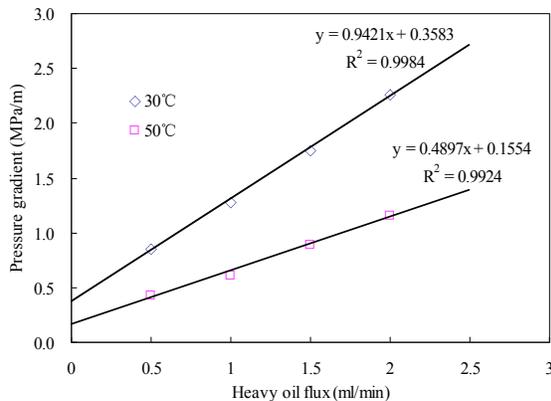


Fig. (2). Threshold pressure gradient vs. heavy oil flux at different temperature.

The results of experiment and calculation were both shown in Fig. (3), which depicted the relationships between dimensionless pressure and dimensionless time at different

temperature. The relative errors of threshold pressure gradient were respectively 0.79% and 2.18% between steady method and transient method. Therefore, the calculated threshold pressure gradients from transient method were reliable as well as the steady state method. However, the transient method only measured the inlet pressure of sand-pack with time after the outlet was opened. Therefore, the transient method need not consider the influence of flow rate on experimental results. Meanwhile, only once experiment was carried out to calculate the threshold pressure gradient of heavy oil. But the steady-state method need measure the relationships between pressure gradients and flow rates, especially, low flow rates of heavy oil in porous media. Therefore, the lower accuracy, which mainly resulted from measuring low flow-rate and small pressure difference, introduced a larger error for threshold pressure gradient in the steady-state method.

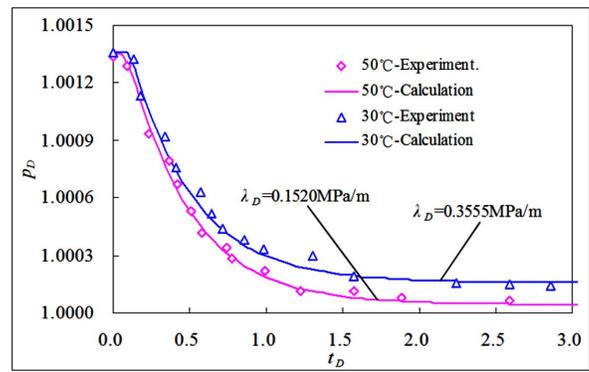


Fig. (3). Comparison of dimensionless pressure between experiments and calculations at different temperature.

3.4. Calculation Method of Threshold Pressure Gradient

The basic parameters, such as sand-pack length, sand-pack diameter, permeability, porosity, heavy oil viscosity, heavy oil density, isothermal compressibility, and a series of threshold pressure gradients were introduced into the equation (15) to calculate the corresponding curves of dimensionless pressure vs. dimensionless time. Then the equation (16) was employed to calculate the corresponding pressure differences along porous media for different threshold pressure gradients. The value of threshold pressure gradient was constantly adjusted until the calculated pressures were matched to the experimental pressures. Therefore, the adjusted value was the real threshold pressure gradient of this heavy oil.

$$\Delta p = \frac{1}{c_o} \text{Ln}(p_D) \tag{16}$$

If the basic parameters were utilized from the former experiments at 30 °C, and the values of threshold pressure gradient were respectively chosen 0 MPa/m, 0.5 MPa/m, 1MPa/m, and 1.5MPa/m, then the curves of pressure differences with dimensionless time could be calculated according to the equations (15) and (16), shown in Fig. (4). It presented the calculated results of the top 100 and the top 200 of the sum expression in the equation (15). The results showed that the lines (the top 100) and the symbols (the top

200) were completely coincident to each other in Fig. (4). Therefore, the results of top 100 in the equation (15) were wholly satisfied with accuracy.

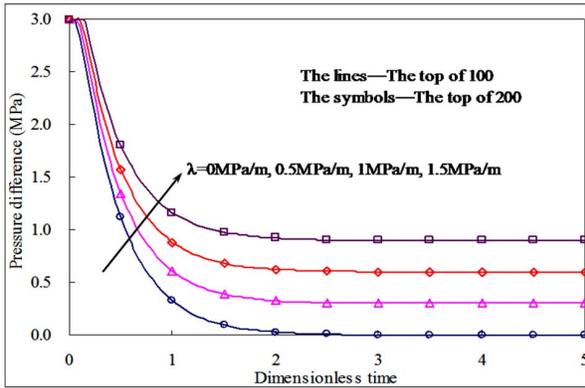


Fig. (4). Comparison of pressure difference between the top of 100 and 200 in equation (9).

3.5. Prediction of Steady Time

When dimensionless time was over 4, pressure distribution gradually tended to be steady in Fig. (4). Therefore, data acquisition time in experiments could be calculated from the following equation:

$$t = 4 \times 1000 \frac{\phi \mu_o c_o L^2}{K} \quad (17)$$

According to the basic parameters in the former experiments, such as $\phi=0.37$, $\mu_o=6110$ mPa·s, $c_o=4.5 \times 10^{-3}$ MPa⁻¹, $L=0.6$ m, $K=2.55$ μm², the time of steady state was 5434.3 s according to the equation (17). In other words, the inlet pressure gradually tended to be steady for about 1.5 hours from opening the outlet of sand-pack.

4. EXPERIMENTAL RESULTS AND DISCUSSION

In our experiments, three heavy oil samples were utilized to study the relationships between oil viscosity and threshold pressure gradient of heavy oil. The permeability were respectively chosen 6730×10^{-3} μm², 2550×10^{-3} μm² and 1600×10^{-3} μm². The viscosity of three heavy oil were respectively chosen 1820 mPa·s on behalf of low viscosity of ordinary heavy oil, 6110 mPa·s on behalf of middle viscosity of ordinary heavy oil and 10715 mPa·s on behalf of super heavy oil. The transient method was employed to calculated threshold pressure gradients of different heavy oils in different permeable porous media.

4.1. Experimental Results

The threshold pressure gradients of the three heavy oils at different temperature were shown from Fig. (5) to Fig. (7). The threshold pressure gradient of ordinary heavy oil with low viscosity gradually decreased as temperature increased in the same permeability of sand-pack as shown in Fig. (5). When temperature increased to a certain value what was called ‘inflection temperature’, the threshold pressure gradient became zero. At the same temperature, the threshold pressure gradient gradually decreased as permeability increased resulting from seepage resistance reduction in

porous media. The results presented that the threshold pressure gradient became zero when it was at 68°C in sand-pack of 1600×10^{-3} μm², at 65 °C in sand-pack of 2500×10^{-3} μm², and at 60 °C in sand-pack of 6730×10^{-3} μm².

The curves about threshold pressure gradient of ordinary heavy oil of 6110 mPa·s at different temperature were shown in Fig. (6). The threshold pressure was only 216.6 kPa at 30 °C in sand-pack of 6730×10^{-3} μm². But the threshold pressure was over 610.0 kPa in sand-pack of 1600×10^{-3} μm². However, threshold pressure gradient largely decreased as temperature increased. When it was at 60 °C, threshold pressure gradient became to be zero in higher permeable porous media. But the threshold pressure gradient gradually disappeared when it was only higher than 70 °C in lower permeable sand-pack of 1600×10^{-3} μm².

Fig. (7) presented the relationships between threshold pressure gradient and temperature for super heavy oil of 10715 mPa·s in different permeable porous media. The threshold pressure was 538.7 kPa, which was far more than ordinary heavy oil, in higher permeable sand-pack of 6730×10^{-3} μm² at 30 °C. The value was even over 960 kPa for super heavy oil in lower permeable sand-pack. Therefore, the seepage resistance of super heavy oil was larger than lower viscosity heavy oil in porous media. But threshold pressure gradient quickly decreased as temperature increased in porous media. When it was over 70 °C, threshold pressure gradient gradually disappeared for super heavy oil in higher permeable sand-pack, that is, its inflection temperature was about 70 °C. But when temperature was close to 80 °C, threshold pressure gradient became to be zero in middle permeable sand-pack of 2500×10^{-3} μm². Meanwhile, threshold pressure gradient gradually tended to be zero when it was over 85 °C in lower permeable sand-pack of 1600×10^{-3} μm². Therefore, super heavy oil presented strongly non-Newtonian characteristics at reservoir temperature. This kind of heavy oil hardly was driven when it was far lower than inflection temperature.

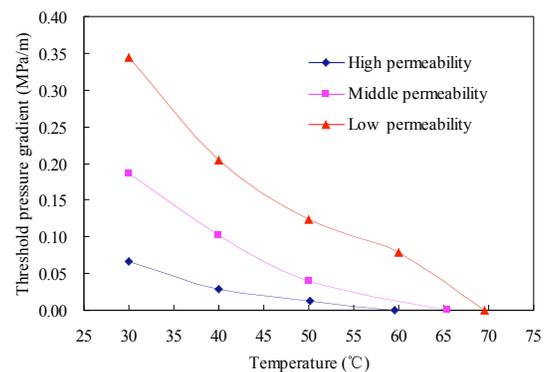


Fig. (5). Threshold pressure gradient vs. temperature of low viscosity of ordinary heavy oil.

4.2. Inflection Temperature

We established the relationships between inflection temperature and heavy-oil viscosity at 50 °C in different permeability, such as 6.73 μm², 2.55 μm² and 1.60 μm² as shown in Fig. (8). The results showed that:

(1) Inflection temperature tended to exponentially increase with viscosity of heavy oil in the same permeable porous media.

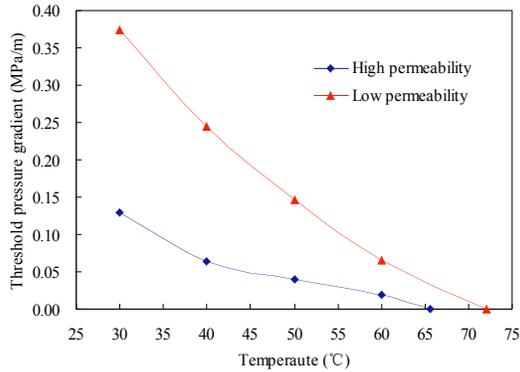


Fig. (6). Threshold pressure gradient vs. temperature of middle viscosity of ordinary heavy oil.

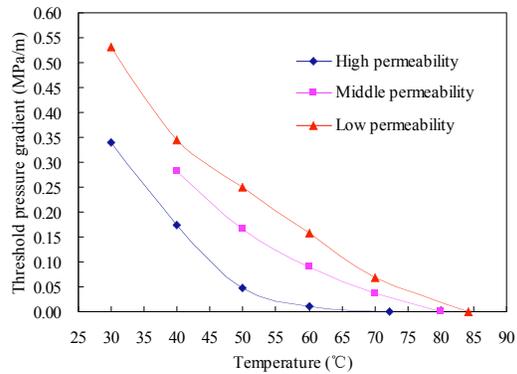


Fig. (7). Threshold pressure gradient vs. temperature of super heavy oil.

(2) Inflection temperature tended to largely decrease as permeability increased for the same oil sample.

The relationships between inflection temperature and viscosity of heavy oil at 50 °C could be expressed by the following equation:

$$T_c = [27.670 - 8.224 \ln(K)] \mu_{0.50}^{0.159} \quad (18)$$

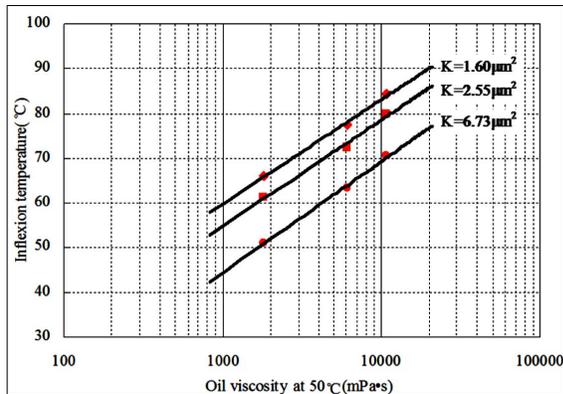


Fig. (8). The fitting curves of inflection temperature and viscosity of heavy oil at 50 °C.

CONCLUSIONS

According to flow characteristics of heavy oil and material balance theory in porous media, an unsteady one

dimension seepage numerical model of heavy oil in porous media was established, which considered threshold pressure gradient of heavy oil during migration. Basis on the theoretical analysis, a new transient experiment method was presented to measure threshold pressure gradient of heavy oil in porous media. This method need not consider flow-rate influence on experimental results, so the transient method was more accurate and need shorter test time than steady state method.

The experimental results show that, on the one hand heavy-oil migration in porous media was influenced by petroleum composition; on the other hand, it was influenced by pore structure of porous media. Inflection temperature of heavy oil presented exponentially increase as viscosity of heavy oil increased in the same permeable porous media.

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CONFLICT OF INTEREST

None declared.

LIST OF SYMBOLS

v	= Seepage velocity of heavy oil, m/s
K	= Absolute permeability of porous media, μm^2
μ_o	= Viscosity of heavy oil, mPa·s
p	= Pressure, Pa
λ	= Threshold pressure gradient of heavy oil, Pa/m
ρ_o	= Density of heavy oil under p , kg/m^3
ϕ	= Porosity of porous media, dimensionless
t	= Time, s
x	= Distance, m
c_o	= Isothermal compressibility of heavy oil, Pa^{-1}
ρ_{osc}	= Density of heavy oil at standard condition, kg/m^3
p_{sc}	= Standard condition pressure, Pa
p_i	= Initial pressure in porous media, Pa
P_D	= Dimensionless pressure
X_D	= Dimensionless distance
t_D	= Dimensionless time
λ_D	= Dimensionless threshold pressure gradient
T	= Temperature, °C
T_c	= Inflection temperature, °C
$\mu_{0.50}$	= Viscosity of heavy oil at 50°C, mPa·s

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